

Feedback Control of Navier-Stokes Equations over Subdomains

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Controlled, 3D, Incompressible Navier–Stokes System

$$\mathbf{w}_t - \nu_0 \Delta \mathbf{w} + (\mathbf{w} \cdot \nabla) \mathbf{w} = -\nabla p + \mathbf{f}(\mathbf{x}, t) + \mathbf{U}$$

$$\operatorname{div} \mathbf{w} = 0, \quad \mathbf{w}|_{\partial\Omega} = 0, \quad \mathbf{w}|_{t=0} = \mathbf{w}_0(\mathbf{x})$$

- $\mathbf{x} = (x_1, x_2, x_3) \in \Omega \subset \mathbb{R}^3, \quad t \in (0, \infty),$
 - Ω – bounded domain
 - $\mathbf{w} = \mathbf{w}(\mathbf{x}, t) = (w_1, w_2, w_3)$ – velocity field
 - p – pressure
 - $\mathbf{w}_t + (\mathbf{w} \cdot \nabla) \mathbf{w}$ – convective derivative,
 - $\nu_0 \sim \frac{1}{\operatorname{Re}}$ – kinetic viscosity, Re – Reynolds #,
 - $\nu_0 \Delta \mathbf{w}$ – diffusion term, “friction,”
 - $\mathbf{w}|_{t=0} = \mathbf{w}_0(\mathbf{x})$ – initial data
 - $\operatorname{div} \mathbf{w} = 0$ – incompressibility,
 - \mathbf{f} – external body force
 - $\mathbf{w}|_{\partial\Omega} = 0$ – “no-slip” boundary condition,
 - \mathbf{U} – control
- ← unknowns

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$$\begin{aligned} \mathbf{w}_t - \nu_0 \Delta \mathbf{w} + (\mathbf{w} \cdot \nabla) \mathbf{w} &= -\nabla p + \mathbf{f}(\mathbf{x}, t) + \mathbf{U} \\ \operatorname{div} \mathbf{w} &= 0, \quad \mathbf{w}|_{\partial\Omega} = 0, \quad \mathbf{w}|_{t=0} = \mathbf{w}_0(\mathbf{x}) \end{aligned}$$

Feedback control over $\Omega_1 \subsetneq \Omega$

$$\begin{aligned} \mathbf{U} = \mathbf{U}(\mathbf{x}, t) &= K\mathbf{w} \\ &= \nu_1 \left(\int_{\Omega_1} |\mathbf{w}_x(\mathbf{x}, t)|^2 d\mathbf{x} \right) \Delta \mathbf{w}(\mathbf{x}, t) \chi_{\Omega_1}(\mathbf{x}) \end{aligned}$$

- $\mathcal{E}_{\Omega_1}(t) = \int_{\Omega_1} |\mathbf{w}_x(\mathbf{x}, t)|^2 d\mathbf{x}$ – energy dissipation rate
- $\mathcal{D}(\mathbf{x}, t) = \chi_{\Omega_1}(\mathbf{x}) \Delta \mathbf{w}(\mathbf{x}, t)$ – friction
- $\chi_{\Omega_1}(\mathbf{x})$ – characteristic function of the set Ω_1
- Output (measurement) = $\mathbf{Y} = (\mathcal{E}_{\Omega_1}(t), \mathcal{D}(\mathbf{x}, t))$
- Input (control) = $\mathbf{U} = K\mathbf{Y} = \nu_1 \mathcal{E}_{\Omega_1}(t) \mathcal{D}(\mathbf{x}, t)$
- $\nu_1 > 0$ – control gain

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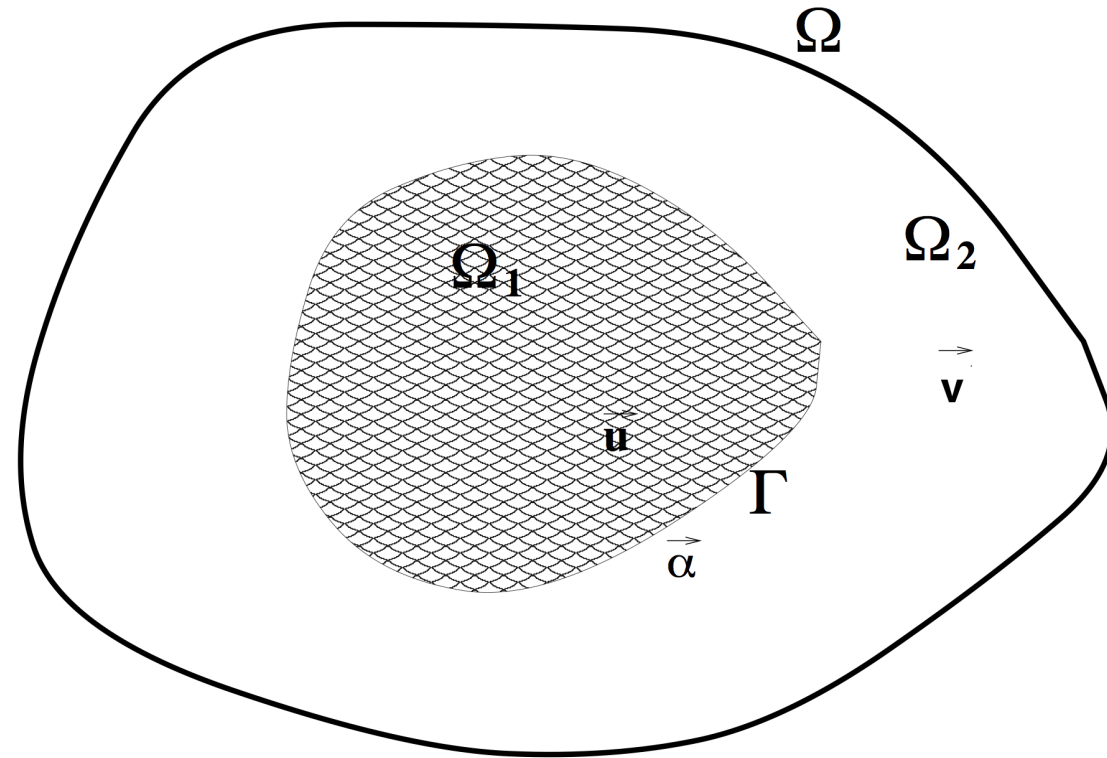
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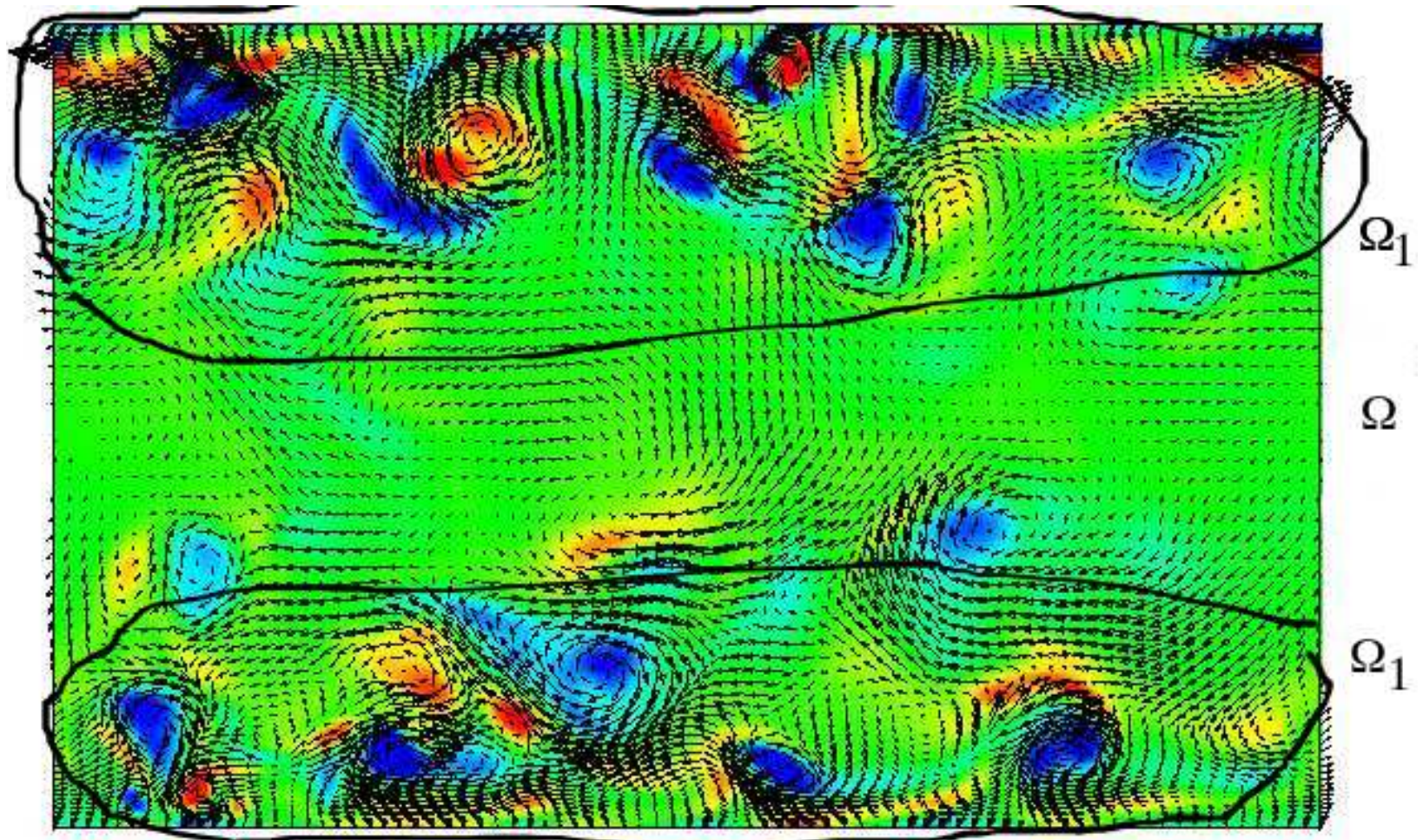
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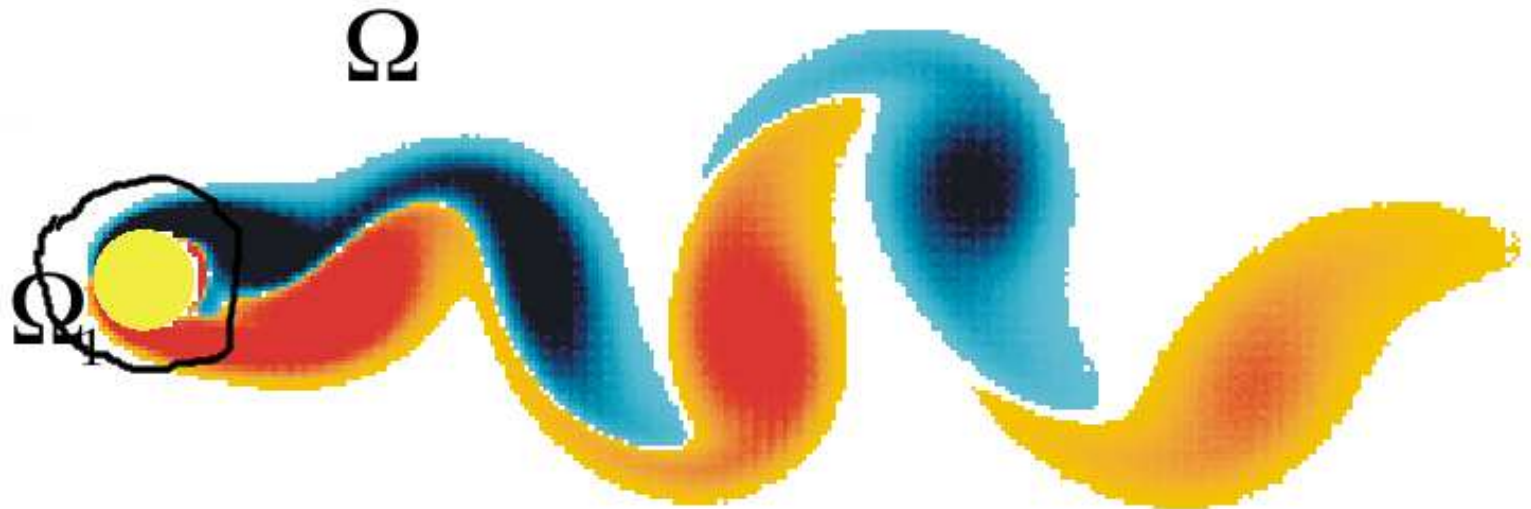
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- In $\Omega_1 \times [0, T]$ (controlled domain)

$$\mathbf{v}_t - \left(\nu_0 + \nu_1 \|\mathbf{v}_x\|^2 \right) \Delta \mathbf{v} + (\mathbf{v} \cdot \nabla) \mathbf{v} = -\nabla p_1 + \mathbf{f}(\mathbf{x}, t)$$

$$\operatorname{div} \mathbf{v} = 0, \text{ in } \Omega_1 \times [0, T], \quad \mathbf{v}|_{\partial\Omega} = 0, \quad \mathbf{v}|_{t=0} = \mathbf{v}_0 \text{ in } \Omega_1$$

- In $\Omega_2 \times [0, T]$, (uncontrolled domain) $\Omega_2 = \Omega \setminus \Omega_1$

$$\mathbf{u}_t - \nu_0 \Delta \mathbf{u} + (\mathbf{u} \cdot \nabla) \mathbf{u} = -\nabla p_2 + \mathbf{f}(\mathbf{x}, t)$$

$$\operatorname{div} \mathbf{u} = 0, \text{ in } \Omega_2 \times [0, T], \quad \mathbf{u}|_{\partial\Omega} = 0, \quad \mathbf{u}|_{t=0} = \mathbf{u}_0 \text{ in } \Omega_2$$

- Interface Conditions on $\Gamma_T = \Gamma \times [0, T]$, ($\Gamma = \partial\Omega_1 \setminus \partial\Omega$)

$$\left(1 + \frac{\nu_1}{\nu_0} \int_{\Omega_1} |\mathbf{v}_x|^2 dx \right) \mathbf{v}_{x_k} \Big|_{\Gamma_T} = \mathbf{u}_{x_k} \Big|_{\Gamma_T}, \quad k = 1, 2, 3$$

$$\mathbf{v} \Big|_{\Gamma_T} = \mathbf{u} \Big|_{\Gamma_T}, \quad p_1 \Big|_{\Gamma_T} = p_2 \Big|_{\Gamma_T}$$

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- spatially nonhomogeneous viscosity coefficient

$$\nu(\mathbf{x}) = \nu_0 + \nu_1 \chi_{\Omega_1}(\mathbf{x}) \left(\int_{\Omega_1} |\mathbf{w}_x|^2 dx \right)$$

- kinetic energy of the fluid $(\frac{1}{2} \|\mathbf{w}\|_{\Omega}^2)$ is a decreasing function of time

$$\frac{1}{2} \frac{d}{dt} \|\mathbf{w}\|_{\Omega}^2 + \nu_0 \|\mathbf{w}_x\|_{\Omega}^2 + \nu_1 \|\mathbf{w}_x\|_{\Omega_1}^4 = 0$$

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O.A. Ladyzhenskaya, “New equations for the description of motion of viscous incompressible fluids and solvability in the large of boundary value problems for them,” *Proc. Steklov Inst. Math.*, 102, (1967), 95–118.

$$\Omega_1 = \Omega$$

- Based on statistical fluid mechanics
- Corresponds to the case of control over the whole domain
-
- Interface condition makes it more complicated
- Domains and the solutions are “glued together”.

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Main Theorem (Assumptions)

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Let \mathbf{f} , \mathbf{w}_0 , and ν_0 satisfy the assumptions

$$\mathbf{f}\Big|_{\Omega_1} \in H^1([0, T], L^2(\Omega_1)) \text{ is arbitrary, and} \quad (1)$$

$$\mathbf{w}_0\Big|_{\Omega_1} \in H^2(\Omega_1) \text{ is arbitrary,} \quad (2)$$

$$\mathbf{f}\Big|_{\Omega \setminus \Omega_1} \in L^2(Q_T'') \text{ for } Q_T'' = \Omega \setminus \Omega_1 \times [0, T], \quad (3)$$

$$\mathbf{w}_0\Big|_{\Omega \setminus \Omega_1} \in H^1(\Omega \setminus \Omega_1), \quad (4)$$

$$\frac{1}{\nu_0} \left(\|\mathbf{w}_0\|_{\Omega \setminus \Omega_1} + \|(\mathbf{w}_0)_x\|_{\Omega \setminus \Omega_1} + \|\mathbf{f}\|_{Q_T''} \right) \leq \rho, \quad (5)$$

where ρ is small enough.

Main Theorem (Statement)

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For any $T > 0$, under the hypotheses (1)–(5) there exists a unique strong solution $\mathbf{w}(\mathbf{x}, t)$ to the controlled Navier–Stokes equations, i.e., there exists a unique vector field $\mathbf{w} \in C([0, T], H^1(\Omega))$, such that $\mathbf{v} = \mathbf{w}|_{\Omega_1} \in C([0, T], H^2(\Omega_1))$, $\mathbf{u} = \mathbf{w}|_{\Omega \setminus \Omega_1} \in C([0, T], H^2(\Omega \setminus \Omega_1))$. More precisely

$$\operatorname{ess\,sup}_{0 \leq t \leq T} \{ \|\mathbf{w}\|_{\Omega}^2 + \|\mathbf{w}_{\mathbf{x}}\|_{\Omega}^2 \} + \int_0^T \left(\|\Delta \mathbf{w}\|_{\Omega}^2 + \|\Delta \mathbf{w}\|_{\Omega \setminus \Omega_1}^2 \right) dt < \infty.$$

The trivial solution of this system is asymptotically stable in the sense that

$$\|\mathbf{w}(\cdot, t)\|_{\Omega} \xrightarrow{t \rightarrow \infty} 0, \text{ when } \|\mathbf{f}(\cdot, t)\|_{\Omega} \xrightarrow{t \rightarrow \infty} 0.$$

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- Consider the equations on the subdomains Ω_1 and $\Omega \setminus \Omega_1$ separately with general boundary conditions
- Solve the Navier–Stokes equations on these subdomains
- Show that there exists boundary condition that satisfies the interface conditions
- Show uniqueness
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- Solve the Navier–Stokes equations on these subdomains
- Show that there exists boundary condition that satisfies the interface conditions
- Show uniqueness
- \vdots
- **Classical energy estimates**

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- Solve the Navier–Stokes equations on these subdomains
- Show that there exists boundary condition that satisfies the interface conditions
- Show uniqueness
- \vdots
- **Classical energy estimates**
- **Contraction mapping with bootstrapping**

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- $\mathcal{N}(\Gamma) = \left\{ \boldsymbol{\alpha} \mid \boldsymbol{\alpha}, \boldsymbol{\alpha}_t \in C\left([0, T], H^{3/2}(\Gamma)\right) \right\}$ (functions space of interface boundary values)
- Subproblem 1. in $Q'_T = \Omega_1 \times [0, T]$ (controlled)

$$\mathbf{v}_t - \nu_0 \Delta \mathbf{v} - \nu_1 \|\mathbf{v}_x\|_{\Omega_1}^2 \Delta \mathbf{v} + v_k v_{x_k} = -\nabla p_1 + \mathbf{f} \quad \text{in } Q'_T$$

$$\operatorname{div} \mathbf{v} = 0 \quad \text{in } Q'_T$$

$$\mathbf{v} = \boldsymbol{\alpha} \quad \text{on } \Gamma_T, \quad \mathbf{v}(\cdot, 0) = \mathbf{v}_0 \quad \text{in } \Omega_1$$

- Subproblem 2. in $Q''_T = (\Omega \setminus \Omega_1) \times [0, T]$ (uncontrolled)

$$\mathbf{u}_t - \nu_0 \Delta \mathbf{u} + u_k \mathbf{u}_{x_k} = -\nabla p_2 + \mathbf{f} \quad \text{in } Q''_T$$

$$\operatorname{div} \mathbf{u} = 0 \quad \text{in } Q''_T$$

$$\mathbf{u} = \boldsymbol{\alpha} \quad \text{on } \Gamma_T, \quad \mathbf{u} = 0 \quad \text{on } \partial\Omega \times [0, T]$$

$$\mathbf{u}(\cdot, 0) = \mathbf{u}_0 \quad \text{in } \Omega \setminus \Omega_1$$

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$$\begin{aligned} \mathbf{v}_t - \nu_0 \Delta \mathbf{v} - \nu_1 \|\mathbf{v}_x\|_{\Omega_1}^2 \Delta \mathbf{v} + v_k \mathbf{v}_{x_k} &= -\nabla p_1 + \mathbf{f} \quad \text{in } Q'_T \\ \operatorname{div} \mathbf{v} &= 0 \quad \text{in } Q'_T \\ \mathbf{v} &= \boldsymbol{\alpha} \quad \text{on } \Gamma_T, \quad \mathbf{v}(\cdot, 0) = \mathbf{v}_0 \quad \text{in } \Omega_1 \end{aligned}$$

$$\begin{aligned} \mathcal{E}_0 : H^{k-\frac{1}{2}}(\Gamma) &\longrightarrow H^k(\Omega), \quad \alpha \longrightarrow \tilde{\alpha} \quad \text{lifting operator} \\ \mathbf{y} &= \mathbf{v} - \tilde{\alpha} \end{aligned}$$

$$\begin{aligned} \mathbf{y}_t - \nu_0 \Delta \mathbf{y} - \nu_1 \left\| (\mathbf{y} + \tilde{\alpha})_x \right\|_{\Omega_1}^2 \Delta (\mathbf{y} + \tilde{\alpha}) + y_k \mathbf{y}_{x_k} \\ + \tilde{\alpha}_k \mathbf{y}_{x_k} + \tilde{\alpha}_{x_k} y_k &= -\nabla p_1 + \mathbf{F}_1 \quad \text{in } Q'_T, \quad (\text{Ctrlld}) \\ \operatorname{div} \mathbf{y} &= 0 \quad \text{in } Q'_T, \\ \mathbf{y} &= 0 \quad \text{on } \Gamma_T, \quad \mathbf{y}(\cdot, 0) = \mathbf{v}_0 - \tilde{\alpha}(\cdot, 0) \quad \text{in } \Omega_1, \end{aligned}$$

$$\mathbf{F}_1 \equiv \mathbf{f} - \tilde{\alpha}_k \tilde{\alpha}_{x_k} + \nu_0 \Delta \tilde{\alpha} - \tilde{\alpha}_t$$

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$$\begin{aligned} & \mathbf{y}_t - \nu_0 \Delta \mathbf{y} - \nu_1 \left\| (\mathbf{y} + \tilde{\boldsymbol{\alpha}})_x \right\|_{\Omega_1}^2 \Delta (\mathbf{y} + \tilde{\boldsymbol{\alpha}}) + y_k \mathbf{y}_{x_k} \\ & + \tilde{\alpha}_k \mathbf{y}_{x_k} + \tilde{\boldsymbol{\alpha}}_{x_k} y_k = -\nabla p_1 + \mathbf{F}_1 \quad \text{in } Q'_T, \quad (\text{Ctrl}d) \\ & \operatorname{div} \mathbf{y} = 0 \quad \text{in } Q'_T, \\ & \mathbf{y} = 0 \quad \text{on } \Gamma_T, \quad \mathbf{y}(\cdot, 0) = \mathbf{v}_0 - \tilde{\boldsymbol{\alpha}}(\cdot, 0) \quad \text{in } \Omega_1, \end{aligned}$$

$$\int_{\Omega_1} (\text{Ctrl}d) \cdot \mathbf{y} d\mathbf{x}$$

$$\begin{aligned} \|\mathbf{y}\|_{\Omega_1}^2 & \leq \int_0^t \left(\frac{2K_{\tilde{\boldsymbol{\alpha}}}^2 K_1}{\nu_1} + 202\nu_1 \|\tilde{\boldsymbol{\alpha}}_x\|_{\Omega_1}^4 + \frac{\|\mathbf{F}_1\|_{\Omega_1}^2 K_1}{\nu_0} \right) ds \\ & + \|\mathbf{y}_0\|_{\Omega_1}^2 \leq C(T, \mathbf{y}_0, \mathbf{f}, \Omega_1, \tilde{\boldsymbol{\alpha}}, \nu_0, \nu_1). \quad (C - E1) \end{aligned}$$

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$$\begin{aligned} \|\mathbf{y}\|_{\Omega_1}^2 &\leq \|\mathbf{y}_0\|_{\Omega_1}^2 + \int_0^t \left(\frac{2K_{\tilde{\alpha}}^2 K_1}{\nu_1} + 202\nu_1 \|\tilde{\alpha}_x\|_{\Omega_1}^4 + \frac{\|\mathbf{F}_1\|_{\Omega_1}^2 K_1}{\nu_0} \right) ds \\ &\leq C(T, \mathbf{y}_0, \mathbf{f}, \Omega_1, \tilde{\alpha}, \nu_0, \nu_1). \quad (C - E1) \Leftarrow \end{aligned}$$

$$\begin{aligned} \frac{d}{dt} \|\mathbf{y}\|_{\Omega_1}^2 + 2\nu_0 \|\mathbf{y}_x\|_{\Omega_1}^2 + \frac{\nu_1}{2} \|\mathbf{y}_x\|_{\Omega_1}^4 + 4\nu_1 \left(\int_{\Omega_1} \langle \mathbf{y}_x, \tilde{\alpha}_x \rangle dx \right)^2 \\ \leq \frac{2K_{\tilde{\alpha}}^2 K_1}{\nu_1} + 202\nu_1 \|\tilde{\alpha}_x\|_{\Omega_1}^4 + \frac{\|\mathbf{F}_1\|_{\Omega_1}^2 K_1}{\nu_0} + \frac{\nu_0 \|\mathbf{y}\|_{\Omega_1}^2}{K_1} \implies \int_0^T \cdot dt \end{aligned}$$

$$\begin{aligned} \|\mathbf{y}\|_{\Omega_1}^2 + 2\nu_0 \int_0^t \|\mathbf{y}_x\|_{\Omega_1}^2 ds + \frac{\nu_1}{2} \int_0^t \|\mathbf{y}_x\|_{\Omega_1}^4 ds + 4\nu_1 \int_0^t \left(\int_{\Omega_1} \langle \mathbf{y}_x, \tilde{\alpha}_x \rangle dx \right)^2 ds \\ \leq C(T, \mathbf{y}_0, \mathbf{f}, \Omega_1, \tilde{\alpha}, \nu_0, \nu_1) \quad (C - E2) \end{aligned}$$

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$$\begin{aligned} \mathbf{y}_t - \nu_0 \Delta \mathbf{y} - \nu_1 \left\| (\mathbf{y} + \tilde{\boldsymbol{\alpha}})_{\mathbf{x}} \right\|_{\Omega_1}^2 \Delta (\mathbf{y} + \tilde{\boldsymbol{\alpha}}) + y_k \mathbf{y}_{x_k} \\ + \tilde{\alpha}_k \mathbf{y}_{x_k} + \tilde{\boldsymbol{\alpha}}_{x_k} y_k = -\nabla p_1 + \mathbf{F}_1 \quad \text{in } Q'_T, \quad (\text{Ctrlld}) \\ \operatorname{div} \mathbf{y} = 0 \quad \text{in } Q'_T, \\ \mathbf{y} = 0 \quad \text{on } \Gamma_T, \quad \mathbf{y}(\cdot, 0) = \mathbf{v}_0 - \tilde{\boldsymbol{\alpha}}(\cdot, 0) \quad \text{in } \Omega_1, \end{aligned}$$

$$\int_{\Omega_1} (\text{Ctrlld})_t \cdot \mathbf{y}_t dx$$

$$\|\mathbf{y}_t\|_{\Omega_1}^2 + \nu_0 \int_0^t \|\mathbf{y}_{xt}\|_{\Omega_1}^2 dt \leq C_4(T, \mathbf{y}_0, \mathbf{f}, \Omega_1, \tilde{\boldsymbol{\alpha}}, \nu_0, \nu_1). \quad (C - E3)$$

also

$$\frac{d}{dt} \|\mathbf{y}_{\mathbf{x}}\|_{\Omega_1}^2 = 2 \int_{\Omega_1} \langle \mathbf{y}_{\mathbf{x}}, \mathbf{y}_{xt} \rangle dx \leq 2 \|\mathbf{y}_{\mathbf{x}}\|_{\Omega_1} \|\mathbf{y}_{xt}\|_{\Omega_1}$$

$$\frac{d}{dt} \|\mathbf{y}_{\mathbf{x}}\|_{\Omega_1} \leq \|\mathbf{y}_{xt}\|_{\Omega_1}$$

$$\|\mathbf{y}_{\mathbf{x}}\|_{\Omega_1} \leq C_5(T, \mathbf{y}_0, \mathbf{f}, \Omega_1, \tilde{\boldsymbol{\alpha}}, \nu_0, \nu_1)$$

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$$\begin{aligned} & \mathbf{y}_t - \nu_0 \Delta \mathbf{y} - \nu_1 \left\| (\mathbf{y} + \tilde{\boldsymbol{\alpha}})_{\mathbf{x}} \right\|_{\Omega_1}^2 \Delta (\mathbf{y} + \tilde{\boldsymbol{\alpha}}) + y_k \mathbf{y}_{x_k} \\ & + \tilde{\boldsymbol{\alpha}}_k \mathbf{y}_{x_k} + \tilde{\boldsymbol{\alpha}}_{x_k} y_k = -\nabla p_1 + \mathbf{F}_1 \quad \text{in } Q'_T, \quad (\text{Ctrlld}) \\ & \operatorname{div} \mathbf{y} = 0 \quad \text{in } Q'_T, \\ & \mathbf{y} = 0 \quad \text{on } \Gamma_T, \quad \mathbf{y}(\cdot, 0) = \mathbf{v}_0 - \tilde{\boldsymbol{\alpha}}(\cdot, 0) \quad \text{in } \Omega_1, \end{aligned}$$

$$\int_{\Omega_1} (\text{Ctrlld}) \cdot \tilde{\Delta} \mathbf{y} \, d\mathbf{x}$$

$$\tilde{\Delta} = P\Delta,$$

$$P : L^2(\Omega_1) \longrightarrow \mathcal{H}(\Omega_1) = \{ \mathbf{y} \in L^2(\Omega_1), \operatorname{div} \mathbf{y} = 0, \mathbf{y} \cdot \boldsymbol{\eta}|_{\Gamma} = 0 \} \quad (\text{Stokes op})$$

$$\| \mathbf{y}_{\mathbf{x}} \|_{\Omega_1}^2 + \nu_0 \int_0^t \left\| \tilde{\Delta} \mathbf{y} \right\|_{\Omega_1}^2 \, ds \leq \| \mathbf{y}_0_{\mathbf{x}} \|_{\Omega_1}^2 + \int_0^t F_4(s) \, ds,$$

$$\nu_0 \int_0^t \left\| \tilde{\Delta} \mathbf{y} \right\|_{\Omega_1}^2 \, ds \leq C(T, \mathbf{y}_0, \mathbf{f}, \Omega_1, \tilde{\boldsymbol{\alpha}}, \nu_0, \nu_1)$$

Existence of Solution on the Controlled Domain

Galerkin's Method

$$\mathbf{y}^m = \sum_{i=1}^m c_{im}(t) \mathbf{a}_i(\mathbf{x}), \quad m = 1, 2, \dots$$

$$\begin{aligned} & \int_{Q'_T} \mathbf{y}_t^m \cdot \phi \, d\mathbf{x} \, dt - \nu_0 \int_{Q'_T} \Delta \mathbf{y}^m \cdot \phi \, d\mathbf{x} \, dt \\ & - \nu_1 \int_{Q'_T} \|(\mathbf{y}^m + \tilde{\boldsymbol{\alpha}})_{\mathbf{x}}\|_{\Omega_1}^2 (\Delta \mathbf{y}^m + \Delta \tilde{\boldsymbol{\alpha}}) \cdot \phi \, d\mathbf{x} \, dt + \int_{Q'_T} y_k^m \mathbf{y}_{x_k}^m \cdot \phi \, d\mathbf{x} \, dt \\ & + \int_{Q'_T} \tilde{\alpha}_k \mathbf{y}_{x_k}^m \cdot \phi \, d\mathbf{x} \, dt + \int_{Q'_T} \tilde{\alpha}_{x_k} \cdot y_k^m \phi \, d\mathbf{x} \, dt = - \int_{Q'_T} \nabla p_1^m \cdot \phi \, d\mathbf{x} \, dt + \int_{Q'_T} \mathbf{F}_1 \cdot \phi \, d\mathbf{x} \, dt \end{aligned}$$

$$\|\mathbf{y}^m\|_{\mathcal{M}_1}^2 \equiv \operatorname{ess\,sup}_{0 \leq t \leq T} \left\{ \|\mathbf{y}^m\|_{\Omega_1}^2 + \|\mathbf{y}_{\mathbf{x}}^m\|_{\Omega_1}^2 + \|\mathbf{y}_t^m\|_{\Omega_1}^2 \right\}$$

$$+ \int_0^T \|\tilde{\Delta} \mathbf{y}^m\|_{\Omega_1}^2 \, dt + \int_0^T \|\mathbf{y}_{\mathbf{x}t}^m\|_{\Omega_1}^2 \, dt \leq C(T, \mathbf{y}_0, \mathbf{f}, \Omega_1, \tilde{\boldsymbol{\alpha}}, \nu_0, \nu_1) < \infty$$

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$$\begin{aligned}\mathbf{u}_t - \nu_0 \Delta \mathbf{u} + u_k \mathbf{u}_{x_k} &= -\nabla p_2 + \mathbf{f} \text{ in } Q_T'' = (\Omega \setminus \Omega_1) \times [0, T], \\ \operatorname{div} \mathbf{u} &= 0 \text{ in } Q_T'', \quad \mathbf{u} = \boldsymbol{\alpha} \text{ on } \Gamma_T, \\ \mathbf{u} &= 0 \text{ on } \partial\Omega \times [0, T], \quad \mathbf{u}(\cdot, 0) = \mathbf{u}_0 \text{ in } \Omega \setminus \Omega_1.\end{aligned}$$

$$\mathbf{y} = \mathbf{v} - \tilde{\boldsymbol{\alpha}}$$

$$\mathbf{y}_t - \nu_0 \Delta \mathbf{y} + y_k \mathbf{y}_{x_k} + \tilde{\boldsymbol{\alpha}}_k \mathbf{y}_{x_k} + \tilde{\boldsymbol{\alpha}}_{x_k} y_k = -\nabla p_2 + \mathbf{F}_1 \text{ in } Q_T''$$

$$\begin{aligned}\operatorname{div} \mathbf{y} &= 0 \text{ in } Q_T'', \quad \mathbf{y} = 0 \text{ on } \partial(\Omega \setminus \Omega_1) \times [0, T], \\ \mathbf{y}(\cdot, 0) &= \mathbf{u}_0 - \tilde{\boldsymbol{\alpha}}(\cdot, 0) \text{ in } \Omega \setminus \Omega_1.\end{aligned}$$

Requirements on the Uncontrolled Domain

We require α , \mathbf{u}_0 , \mathbf{f} to be small enough, and ν_0 to be large enough, so that

$$\begin{aligned} \nu_0 - 2K_{\tilde{\alpha}} \sqrt{K_1} &\geq 0 \\ \nu_0^2 K_1 - 4K_{\tilde{\alpha}}^2 - 4C^2 \|\tilde{\alpha}_x\|_{\Omega \setminus \Omega_1} \|\Delta \tilde{\alpha}\|_{\Omega \setminus \Omega_1} &> 0, \end{aligned}$$

$$\begin{aligned} &\|\mathbf{u}_{0x} - \tilde{\alpha}_x(\cdot, 0)\|_{\Omega \setminus \Omega_1}^2 + \frac{4}{\nu_0} \int_0^T \|\mathbf{F}_1\|_{\Omega \setminus \Omega_1}^2 dt \\ &\leq \frac{\nu_0^2}{6\sqrt{6}C^2} \sqrt{K_1 - \frac{4K_{\tilde{\alpha}}^2}{\nu_0^2} - \frac{4C^2 \|\tilde{\alpha}_x\|_{\Omega \setminus \Omega_1} \|\Delta \tilde{\alpha}\|_{\Omega \setminus \Omega_1}}{\nu_0^2}} \end{aligned}$$

where

$$\mathbf{F}_1 \equiv \mathbf{f} - \tilde{\alpha}_k \tilde{\alpha}_{x_k} + \nu_0 \Delta \tilde{\alpha} - \tilde{\alpha}_t,$$

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Using a priori estimates and Galerkin's method:

There exists a unique solution $\mathbf{u}(x, t)$ to the Navier–Stokes equations on the uncontrolled domain $\Omega \setminus \Omega_1$, which is in the class \mathcal{M}_2 of functions having finite norm

$$\|\mathbf{u}\|_{\mathcal{M}_2}^2 \equiv \operatorname{ess\,sup}_{0 \leq t \leq T} \{ \|\mathbf{u}\|_{\Omega \setminus \Omega_1}^2 + \|\mathbf{u}_x\|_{\Omega \setminus \Omega_1}^2 \} + \int_0^T \|\Delta \mathbf{u}\|_{\Omega \setminus \Omega_1}^2 dt.$$

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Combining the Two Subproblems

$$\mathbf{w}(\mathbf{x}, t) = \begin{cases} \mathbf{v}(\mathbf{x}, t) & \text{in } Q'_T \\ \mathbf{u}(\mathbf{x}, t) & \text{in } Q''_T \end{cases}$$

- $\mathbf{u}|_{\Gamma_T} = \mathbf{v}|_{\Gamma_T} = \boldsymbol{\alpha} \in \mathcal{N}(\Gamma)$
- There exists a unique $\boldsymbol{\alpha}$ such that

$$\left(1 + \frac{\nu_1}{\nu_0} \int_{\Omega_1} |\mathbf{v}_x|^2 dx \right) \mathbf{v}_{x_k} \Big|_{\Gamma_T} = \mathbf{u}_{x_k} \Big|_{\Gamma_T}, \quad k = 1, 2, 3$$

(Fixed point problem)

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D-N Mappings

Define the mappings

$$\Phi, \Psi : \mathcal{N}(\Gamma) \longrightarrow C([0, T], H^{1/2}(\Gamma))$$

$$\alpha \xrightarrow{\Psi} \Psi(\alpha) \equiv \left(1 + \frac{\nu_1}{\nu_0} \int_{\Omega_1} |\mathbf{v}_x|^2 dx \right) \mathbf{v}_x|_{\Gamma}$$



$$\alpha \xrightarrow{\Phi} \Phi(\alpha) \equiv \mathbf{u}_x|_{\Gamma}$$



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Interface Condition

$$\Phi(\alpha) = \Psi(\alpha)$$

or equivalently (fixed point formulation)

$$\mathcal{K}(\alpha) \equiv \Phi^{-1} \circ \Psi(\alpha) = \alpha$$

$\mathcal{K} = \Phi^{-1} \circ \Psi$ is a contraction mapping
on sufficiently small time interval $[0, T_0]$.

1. Ψ is Lipschitz continuous:

$$\|\Psi(\alpha_1) - \Psi(\alpha_2)\|_{C([0, T], H^{1/2}(\Gamma))} \leq C \|\alpha_1 - \alpha_2\|_{\mathcal{N}(\Gamma)}$$

- embedding theorems
- trace theorem
- extension theorem

Finding Fixed Point (Continued)

2. Φ^{-1} exists and is Lipschitz continuous.

$$\|\alpha_1 - \alpha_2\|_{\mathcal{N}(\Gamma)} \leq C(T) \|\Phi(\alpha_1) - \Phi(\alpha_2)\|_{C([0,T], H^{1/2}(\Gamma))}$$
$$C(T) \xrightarrow{T \rightarrow 0} 0$$

- (a) Φ is one-to-one, hence invertible
- (b) Φ^{-1} is a contraction mapping for small T_0

- embedding theorems
- trace theorem
- extension theorem
- generalized Poincaré inequality (Lawson)
- pressure estimates

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- $\mathcal{K} = \Phi^{-1} \circ \Psi$ is a contraction mapping for sufficiently small time interval $[0, T_0]$
- $C(T)$ does not depend on the data
- Time steps of length T_0 cover an arbitrary interval $[0, T]$

One of many estimates

$$\begin{aligned} \|\mathbf{w}(\cdot, t)\|_{\Omega} &\leq \|\mathbf{w}(\cdot, 0)\|_{\Omega} e^{-\nu_0 K_1 t} \\ &+ \int_0^t \|\mathbf{f}(\mathbf{x}, s)\|_{\Omega} e^{-\nu_0 K_1 (t-s)} ds \end{aligned}$$

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Joint Subproblems

D-N Mappings

- Promising preliminary results in a 2D channel with control in the boundary layer.
- Plans for 3D problems using GPUs.